

## Maxwell 方程式

9 24 04 2019 基礎

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

$$\nabla \times \mathbf{E} = - \frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{j} + \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t}$$

$\mathbf{E}$  : 電場

$\mathbf{B}$  : 磁束密度

$\rho$  : 電荷密度

$\mu_0$  : 透磁率

$\epsilon_0$  : 誘電率

$\mathbf{E}(r, t)$

$\mathbf{B}(r, t)$

$\rho(r)$

$\mathbf{j}(r)$

$\phi$   
場所 (r, t)

u3!!

# 0-2 数字のまとめ

No

Date

6

0-2-1 ベクトル 内積, 外積  
スカラー積      ベクトル積

$$\left\{ \begin{aligned} \mathbf{a} \cdot \mathbf{b} &= a_x b_x + a_y b_y + a_z b_z \end{aligned} \right.$$

$$\mathbf{a} \times \mathbf{b} = (a_y b_z - a_z b_y) \mathbf{e}_x + (a_z b_x - a_x b_z) \mathbf{e}_y + (a_x b_y - a_y b_x) \mathbf{e}_z = \begin{vmatrix} \mathbf{e}_x & \mathbf{e}_y & \mathbf{e}_z \\ a_x & a_y & a_z \\ b_x & b_y & b_z \end{vmatrix}$$

$$\mathbf{a} \times \mathbf{a} = 0$$

0-2-2 微分

$$\nabla \equiv \left( \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right)$$

$\frac{\partial}{\partial x}$  : 偏微分

$$\boxed{\frac{\partial f}{\partial x}}$$

yz z と x の関係

x のみで微分する

◎ 合成微分  $\phi(x, y, z, t)$  のとき

$$\begin{aligned} \frac{d\phi}{dt} &= \frac{\partial\phi}{\partial x} \frac{dx}{dt} + \frac{\partial\phi}{\partial y} \frac{dy}{dt} + \frac{\partial\phi}{\partial z} \frac{dz}{dt} + \frac{\partial\phi}{\partial t} \\ &= (\nabla\phi) \cdot \dot{\mathbf{r}} + \frac{\partial\phi}{\partial t} \end{aligned}$$

ただし  $\nabla\phi \equiv \left( \frac{\partial\phi}{\partial x}, \frac{\partial\phi}{\partial y}, \frac{\partial\phi}{\partial z} \right)$

◎ divergence (発散)

$$\begin{aligned} \nabla \cdot \mathbf{E} &= \frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} \\ &\equiv \text{div } \mathbf{E} \quad \text{と書ける} \end{aligned}$$

物理的な意味:

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0} \quad (\text{Gauss の法則})$$



$\mathbf{E}$  から  $\rho$  による  $\frac{1}{r^2}$  の  $\frac{1}{r^2}$

これは

$\frac{1}{r^2}$  の「源」が  $\rho$  である

③ 微分  $\frac{\partial \phi}{\partial x} \Delta x$   $\approx \frac{\partial \phi}{\partial x} \Delta x$

$$(1) \phi(x(t)) \quad a \neq 0$$

$$\Delta \phi \equiv \phi(x(t) + \Delta x) - \phi(x(t))$$

$$= \frac{\Delta \phi}{\Delta x} \cdot \Delta x = \frac{\partial \phi}{\partial x} \cdot \Delta x$$

$$\frac{d\phi}{dt} = \lim_{\Delta t \rightarrow 0} \frac{\Delta \phi}{\Delta t} = \frac{\partial \phi}{\partial x} \frac{\Delta x}{\Delta t} = \frac{\partial \phi}{\partial x} \frac{dx}{dt}$$

$$(2) \phi(x(t), y(t)) \quad a \neq 0$$

$$\Delta \phi \equiv \phi(x(t) + \Delta x, y(t) + \Delta y) - \phi(x, y)$$

$$= \phi(x + \Delta x, y + \Delta y) - \phi(x, y + \Delta y)$$

$$+ \phi(x, y + \Delta y) - \phi(x, y)$$

$$= \frac{\partial \phi}{\partial x} \cdot \Delta x + \frac{\partial \phi}{\partial y} \cdot \Delta y$$

$$\therefore \begin{cases} \frac{\partial \phi}{\partial x} \equiv \frac{\phi(x + \Delta x, y) - \phi(x, y)}{\Delta x} \\ \frac{\partial \phi}{\partial y} \equiv \frac{\phi(x, y + \Delta y) - \phi(x, y)}{\Delta y} \end{cases}$$

$$\therefore \frac{d\phi}{dt} = \frac{\partial \phi}{\partial x} \frac{dx}{dt} + \frac{\partial \phi}{\partial y} \frac{dy}{dt} //$$

• gradient (電位  $\phi$ )

$$\nabla \phi = \left( \frac{\partial \phi}{\partial x}, \frac{\partial \phi}{\partial y}, \frac{\partial \phi}{\partial z} \right)$$

$$\boxed{\mathbf{E} = -\nabla \phi} \quad \text{in } \mathbb{C}$$



$\phi$ : 電位 (電位)

(スカラーポテンシャル)

なぜ? どのようにして?

$$\boxed{\nabla \times \mathbf{E} = 0}$$

Maxwell eq.

9.1 (2) の結果

$$\mathbf{E} = -\nabla \phi \quad \Rightarrow$$

$$\nabla \times \mathbf{E} = -\nabla \times \nabla \phi = 0$$

これは電場の性質

● No charge ( $\rho=0$ )

$$\nabla \cdot \vec{E} = 0$$

$$\vec{E} = -\nabla\phi \quad \text{in } \mathbb{R}^3$$

$$\boxed{\nabla^2\phi = 0}$$

↓

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \phi = 0$$

একটা সমাধান কি আছে?

Ans:

$$\boxed{\phi = \frac{\alpha}{r}}$$

( $\alpha$ : constant)

( $r \neq 0$  only)

সম্পন্ন করে দেখাও!!

(\*)

$$\frac{\partial}{\partial x} \frac{1}{r} = \left( \frac{\partial}{\partial r} \frac{1}{r} \right) \frac{\partial r}{\partial x} = -\frac{1}{r^2} \frac{x}{r} = -\frac{x}{r^3}$$

$$\frac{\partial^2}{\partial x^2} \frac{1}{r} = -\frac{\partial}{\partial x} \frac{x}{r^3} = -\frac{1}{r^3} - \frac{\partial}{\partial x} \frac{1}{r^3}$$

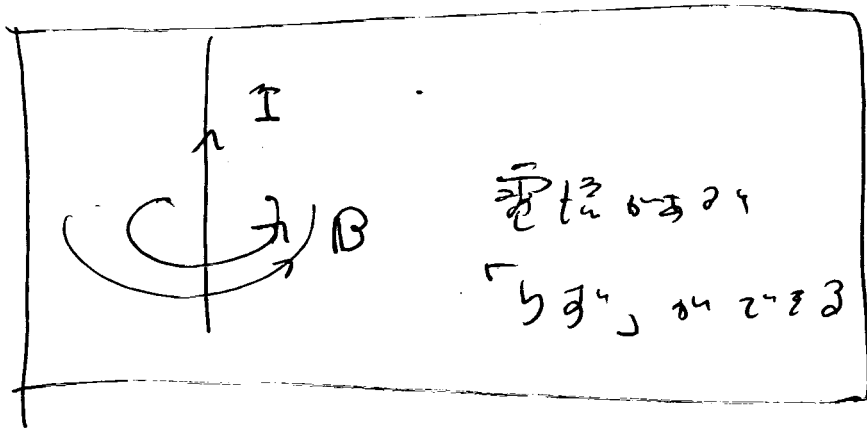
$$= -\frac{1}{r^3} + \frac{3x^2}{r^5}$$

$$\therefore \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \frac{1}{r} = \left( -\frac{3}{r^3} + \frac{3(x^2+y^2+z^2)}{r^5} \right) = 0$$

১০০ (৮) এর সমাধান

② rotation :  $\nabla \times \mathbf{B} = \mu_0 \mathbf{i}$

$$\nabla \times \mathbf{B} = \begin{vmatrix} \mathbf{e}_1 & \mathbf{e}_2 & \mathbf{e}_3 \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ B_x & B_y & B_z \end{vmatrix} \equiv \text{rot } \mathbf{B}$$



# 0-2-3 積分 (Integral)

**積分**: 公式は絶対12算22

少くとも公式が12212

0, 243の12, 24322!!

## 1. 線積分

線1216, 2積分23

積分は1次元

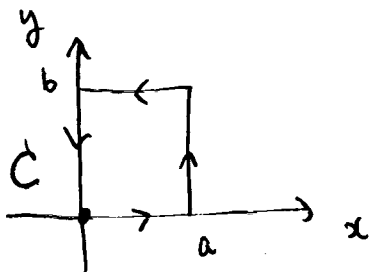
但し3次元空間2行3

$$I = \int_C A \cdot ds$$

A 2115  $\frac{d}{t}$  2 C 2115 線1216, 2  
S積分23

(Example) ①

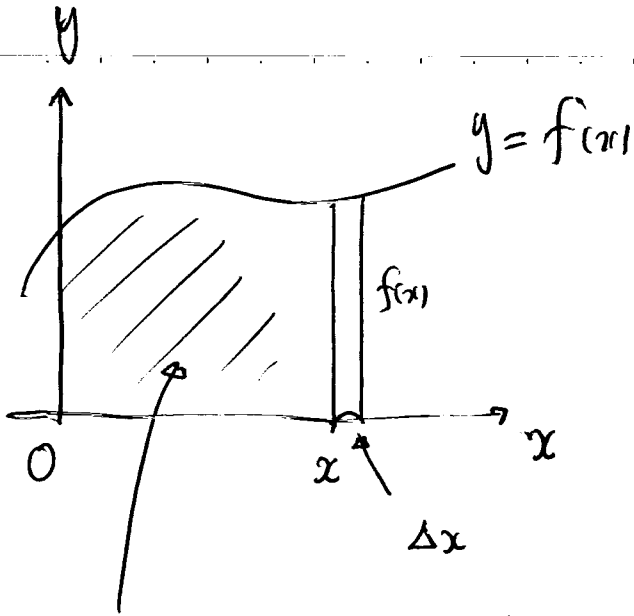
$$A = (x^2 + y^2, xy, 0)$$



$$ds = (dx, dy, 0)$$



10'



$$S(x) = \int_0^x f(x) dx$$

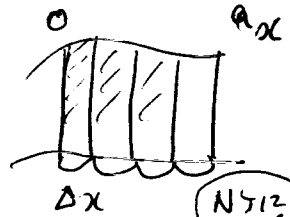
$$S(x + \Delta x) = S(x) + \Delta x \cdot f(x)$$

$$\therefore f(x) = \frac{S(x + \Delta x) - S(x)}{\Delta x}$$

$\Delta x \rightarrow 0$  等价于

$$f(x) = \frac{dS}{dx}$$

积分应用



$$\sum_i f(x_i) \Delta x$$

$$\Delta x = \frac{x}{N}$$

$$x_i = \Delta x \cdot i$$

$$\sum_i f(x_i) \Delta x \quad (N \rightarrow \infty)$$

$$\equiv \int_0^x f(x) dx = S(x)$$

ε 的 ε

## 2. 面積積分

$$I = \int A \cdot dS$$

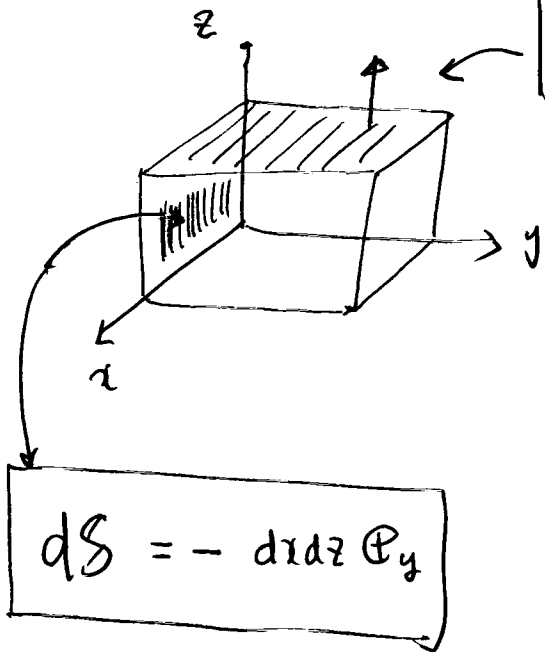
$$dS = m \cdot dS$$



↑ 微小面積  $dS$  の  
その面の法線方向の向きをベクトル

ゆき外向き

(a) 直方体

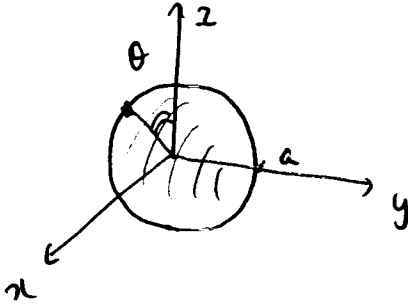


$dS = dx dy \hat{e}_z$

$dS = - dx dz \hat{e}_y$

体積 12212  
常に外向き

z方向定義

(b) 球面上 (半径  $a$ )

$$dS = a^2 \sin\theta d\theta d\varphi \mathbf{e}_r$$

$$\int A \cdot dS$$

$$A = A(r)$$

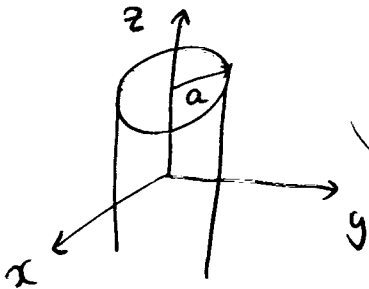
ε q 3

$$= \int A_r \cdot a^2 \sin\theta d\theta d\varphi$$

$$= A_r \cdot a^2 \int_0^\pi \sin\theta d\theta \int_0^{2\pi} d\varphi$$

$$= 4\pi a^2 A_r \quad //$$

(c) 圆柱上



$$\int A \cdot dS$$

$$A = A(r)$$

ε q 3

$$= \int A_r \cdot a d\theta dz$$

$$= 2\pi a \int A_r \cdot dz$$

$$dS = a d\theta dz \cdot \mathbf{e}_r$$

## 3. 体積要素

$$\int d^3r \equiv \int dx dy dz$$

極座標

$$dx dy dz = \frac{\partial(x, y, z)}{\partial(r, \theta, \phi)} dr d\theta d\phi$$

$$\frac{\partial(x, y, z)}{\partial(r, \theta, \phi)} = \begin{vmatrix} \frac{\partial x}{\partial r} & \frac{\partial y}{\partial r} & \frac{\partial z}{\partial r} \\ \frac{\partial x}{\partial \theta} & \frac{\partial y}{\partial \theta} & \frac{\partial z}{\partial \theta} \\ \frac{\partial x}{\partial \phi} & \frac{\partial y}{\partial \phi} & \frac{\partial z}{\partial \phi} \end{vmatrix}$$

$$= \begin{vmatrix} \sin\theta \cos\phi & \sin\theta \sin\phi & \cos\theta \\ r \cos\theta \cos\phi & r \cos\theta \sin\phi & -r \sin\theta \\ -r \sin\theta \sin\phi & r \sin\theta \cos\phi & 0 \end{vmatrix}$$

$$= r^2 \sin\theta$$

即

$$\begin{cases} x = r \sin\theta \cos\phi \\ y = r \sin\theta \sin\phi \\ z = r \cos\theta \end{cases}$$

$$\int d^3r = \int r^2 \sin\theta dr d\theta d\phi$$

$$= \int r^2 dr \int_0^\pi \sin\theta d\theta \int_0^{2\pi} d\phi$$

No. \_\_\_\_\_  
Date \_\_\_\_\_

# 0-3 Gauss の定理 & Stokes の定理

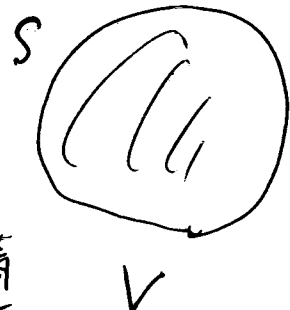
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① Gauss : 
$$\int_V \operatorname{div} \mathbf{A} \cdot d^3r = \int_S \mathbf{A} \cdot d\mathbf{S} = \int_S \mathbf{A} \cdot \mathbf{n} \cdot dS$$

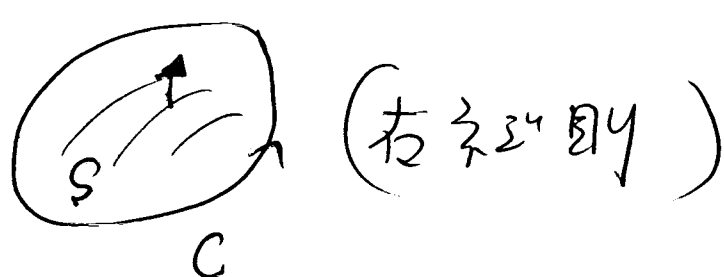
② Stokes : 
$$\int_S \nabla \times \mathbf{A} \cdot d\mathbf{S} = \int_C \mathbf{A} \cdot d\mathbf{r}$$

{ V : 有限の体積

{ S : その体積を包む曲面



S : 曲线 C を包む曲面

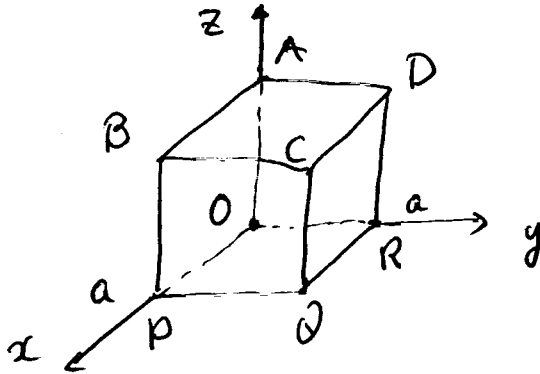


(右ネジ則)

# ① Gauss の定理の導出

[ 3つの Step で示す ]

1. Step 1 | ① の場合の導出を示す



$$\int_{\text{体積}} \operatorname{div} \mathbf{A} \cdot d^3r = \int \frac{\partial A_x}{\partial x} dx dy dz + \int \frac{\partial A_y}{\partial y} dx dy dz + \int \frac{\partial A_z}{\partial z} dx dy dz$$

$$= \int \left[ A_x(x, y, z) \right]_{z=0}^{z=a} dy dz + \int \left[ A_y(x, y, z) \right]_{y=0}^{y=a} dx dz + \int \left[ A_z(x, y, z) \right]_{z=0}^a dx dy$$

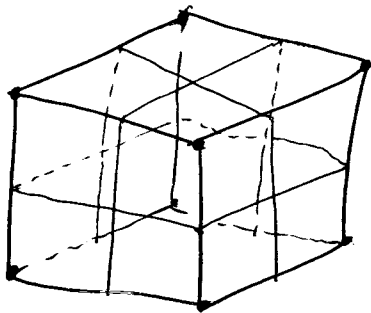
$$= \int (A_x(a, y, z) - A_x(0, y, z)) dy dz + \int (A_y(x, a, z) - A_y(x, 0, z)) dx dz + \int (A_z(x, y, a) - A_z(x, y, 0)) dx dy$$

$$= \int_{\text{面 PQCB}} A_x dydz - \int_{\text{面 ORDA}} A_x dydz + \dots$$

$$= \int_{\text{全表面}} A \cdot n dS$$

$n$ : 法線方向の単位ベクトル  
 経路の3方向外向

[Step 2] 8つの立方体の分割

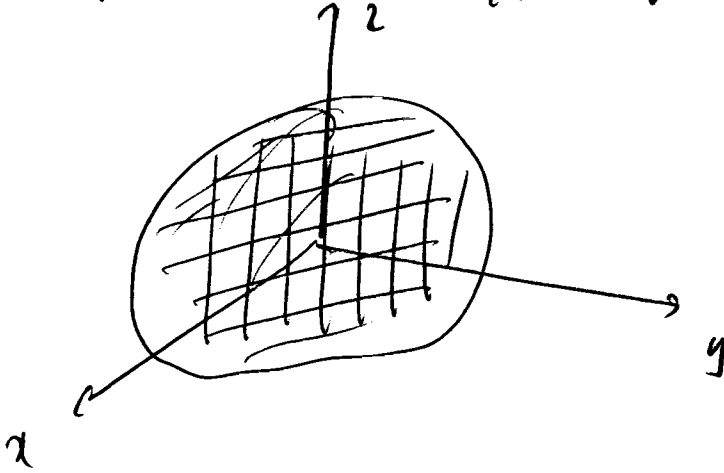


面が重なり、  
 2:3は逆向き  
 2:1は消滅

$$\int_{\text{全体積}} \text{div } A d^3x = \sum_i \int_{[i \text{ 立方体}]} \text{div } A d^3x$$

$$= \sum_i \int_{\substack{i \text{-立方体} \\ (\alpha \text{表面})}} A \cdot n dS = \int_{\substack{\text{全体の} \\ \text{表面}}} A \cdot n dS$$

[Step 3] 任意の体積  $V$



[有限個の立体の分割]

$$\int \operatorname{div} A \, dV = \sum_i \int \operatorname{div} A \, dV_i$$

(i 立体の体積)

$$= \sum_i \int A \cdot n \, dS$$

(↑ 立体の表面積)

電場の発散は閉路の

$$= \int A \cdot n \, dS$$

(全体  
の表面積)

分割  $\epsilon \approx \text{PG} \approx \text{EMC} \approx \text{MS}$

→  $V$



$$\int \operatorname{div} A \, d^3r = \int A \cdot dS$$

【法線】  $m$  :



$$|m| = 1$$

面は垂直に外向  
常に外向の方向

【接線】  $n$  :



$$|n| = 1$$

面は平行

面は

微小面積の意味