

5-4 Separable Interaction

No.

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The Lippman-Schwinger equation for the T -matrix can be written as

$$T = V + V G T$$

More explicitly, we write

$$\begin{aligned} \langle k' | T | k \rangle &= \langle k' | V | k \rangle \\ &+ \sum_{k''} \langle k' | V | k'' \rangle \langle k'' | \frac{1}{E - H_0} | k'' \rangle \\ &\times \langle k'' | T | k \rangle \end{aligned}$$

Here, we define the Green's function in momentum space as

$$G_E(k) \equiv \frac{1}{E - \frac{k^2}{2m} + i\epsilon}$$

Thus, we find

$$\langle k' | T | k \rangle = \langle k' | V | k \rangle + \sum_{k''} \langle k' | V | k'' \rangle G_E(k'') \langle k'' | T | k \rangle$$

- Separable interaction (potential)

We take

$$\langle k' | V | k \rangle = \lambda g(k') g(k)$$

$$\begin{cases} \lambda : \text{some constant} \\ g(k) : \text{vertex} \end{cases}$$

Therefore, we obtain

$$\begin{aligned} \langle k' | T | k \rangle &= \lambda g(k') g(k) + \lambda g(k') \times \\ &\int \frac{d^3 k''}{(2\pi)^3} g(k'') Q_E(k'') \langle k'' | T | k \rangle \end{aligned}$$

We further assume that

$\langle k' | T | k \rangle$ should be written as

$$\langle k' | T | k \rangle = \beta(E) g(k') g(k)$$

In this case, the equation for the T-matrix becomes

$$\beta(E) g(k') g(k) = \lambda g(k') g(k) + \lambda g(k') \beta(E) g(k) \\ \times \int \frac{d^3 k''}{(2\pi)^3} [g(k'')]^2 G_E(k'')$$

Therefore, we find

$$\beta(E) = \lambda + \lambda \beta(E) \int \frac{d^3 k''}{(2\pi)^3} [g(k'')]^2 G_E(k'')$$

We have define

$$\left\{ \begin{array}{l} \bar{F}(E) = \int \frac{d^3 k''}{(2\pi)^3} [g(k'')]^2 G_E(k'') \end{array} \right.$$

$$E : \text{incident energy} \quad \bar{E} = \frac{k^2}{2m}$$

\therefore

$$\beta(E) = \frac{\lambda}{1 - \lambda \bar{F}(E)}$$

Since the T -matrix is assumed to be

$$\langle k' | T | k \rangle = \beta(E) g(k') g(k)$$

we find

$$\langle k' | T | k \rangle = \frac{\lambda}{1 - \lambda F(E)} g(k') g(k)$$

This is the exact solution!

- Note: The separable interaction is not realistic, therefore, a toy model.